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# Classification of low energy sign-changing solutions of an almost critical problem <sup>\*</sup>

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**Abstract.** In this paper we make the analysis of the blow-up of low energy sign-changing solutions of a semi-linear elliptic problem involving nearly critical exponent. Our results allow to classify these solutions according to the concentration speeds of the positive and negative part and, in high dimensions, lead to complete classification of them. Additional qualitative results, such as symmetry or location of the concentration points are obtained when the domain is a ball.

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**Key words:** Blow-up analysis, sign-changing solutions, nodal domains, critical exponent.

## 1 Introduction

In this paper we consider the following semi-linear elliptic problem with subcritical nonlinearity:

$$(1) \quad \begin{cases} -\Delta u = |u|^{2^*-2-\varepsilon}u & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

where  $\Omega$  is a smooth bounded domain in  $\mathbb{R}^n$ ,  $n \geq 3$ ,  $\varepsilon$  is a positive real parameter and  $2^* = 2n/(n-2)$  is the critical Sobolev exponent for the embedding of  $H_0^1(\Omega)$  into  $L^{2^*}(\Omega)$ .

Problem (1) is related to the limiting problem (when  $\varepsilon = 0$ ) which exhibits a lack of compactness and gives rise to solutions of (1) which blow up as  $\varepsilon \rightarrow 0$ .

In the last decades there have been many works devoted to the study of positive solutions of problem (1). In sharp contrast to this, very little study has been made concerning the sign-changing solutions. For details one can see [5] and the references therein.

The existence of sign-changing solutions of (1) for any  $\varepsilon \in (0, p-1)$  has been proved in [4], [6] and [11]. On the other hand, when  $\varepsilon = 0$ , problem (1) becomes delicate. Pohozaev showed in [16] that if  $\Omega$  is starshaped, problem (1) (with  $\varepsilon = 0$ ) has no solutions whereas Clapp and

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Weth proved in [12] that problem (1) (with  $\varepsilon = 0$ ) has a solution on domains with small holes and on some contractible domains with an involution symmetry.

In view of this qualitative change in the situation when  $\varepsilon = 0$ , it is interesting to study the following question: what happens to the solutions of (1) as  $\varepsilon \rightarrow 0$ ? In this paper, we are mainly interested in the study of the behavior of low energy sign-changing solutions of (1). The first part of this paper is devoted to analyze the asymptotic behavior, as  $\varepsilon \rightarrow 0$ , of sign-changing solutions of (1) whose energy converges to the value  $2S^{n/2}$ ,  $S$  being the Sobolev constant for the embedding of  $H_0^1(\Omega)$  into  $L^{2^*}(\Omega)$ . We prove that these solutions blow up at two points which may coincide and which are the limit of the concentration points  $a_{\varepsilon,1}$  and  $a_{\varepsilon,2}$  of the positive and negative part of the solutions. Moreover, we make a precise study of the location of these blow up points. More precisely, we have

**Theorem 1.1** *Let  $n \geq 3$  and let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) which satisfies*

$$\|u_\varepsilon\|^2 := \int_{\Omega} |\nabla u_\varepsilon|^2 \rightarrow 2S^{n/2} \quad \text{as } \varepsilon \rightarrow 0. \quad (1.1)$$

*Then, the set  $\Omega \setminus \{x \in \Omega \mid u_\varepsilon(x) = 0\}$  has exactly two connected components.*

*Furthermore, there exist two points  $a_{\varepsilon,1}$ ,  $a_{\varepsilon,2}$  in  $\Omega$  (one of them can be chosen to be the global maximum point of  $|u_\varepsilon|$ ) and there exist two positive reals  $\mu_{\varepsilon,1}$ ,  $\mu_{\varepsilon,2}$  such that*

$$\|u_\varepsilon - P\delta_{(a_{\varepsilon,1}, \mu_{\varepsilon,1})} + P\delta_{(a_{\varepsilon,2}, \mu_{\varepsilon,2})}\| \rightarrow 0, \quad \text{as } \varepsilon \rightarrow 0, \quad (1.2)$$

$$\mu_{\varepsilon,i} d(a_{\varepsilon,i}, \partial\Omega) \rightarrow +\infty, \quad |\mu_{\varepsilon,i}|^\varepsilon \rightarrow 1 \quad \text{as } \varepsilon \rightarrow 0, \quad \text{for } i \in \{1, 2\}, \quad (1.3)$$

where  $P\delta_{(a,\mu)}$  denotes the projection of  $\delta_{(a,\mu)}$  on  $H_0^1(\Omega)$ , that is,

$$\Delta P\delta_{(a,\mu)} = \Delta\delta_{(a,\mu)} \text{ in } \Omega, \quad P\delta_{(a,\mu)} = 0 \text{ on } \partial\Omega, \quad \text{and } \delta_{(a,\mu)}(x) = \frac{\beta_n \mu^{(n-2)/2}}{(1 + \mu^2|x - a|^2)^{(n-2)/2}}.$$

Here  $\beta_n$  is a constant chosen so that  $-\Delta\delta_{(a,\mu)} = \delta_{(a,\mu)}^{(n+2)/(n-2)}$ , ( $\beta_n = (n(n-2))^{(n-2)/4}$ ).

Note that, for the supercritical case (that is for  $\varepsilon < 0$ ), a recent result [9] shows that there is no sign-changing solution  $u_\varepsilon$  with low energy which satisfies (1.2) and (1.3).

Now, our aim is to give a complete classification of the solutions satisfying (1.1). To this aim, we divide this kind of solutions into two categories: the ones for which the positive and negative part blow up with the same rate (hypothesis (1.5) below) and the ones for which these rates are not comparable (hypothesis (1.9) below). In the first case we are able to prove that the concentration points of the positive and negative part of a solution of this type are distinct and away from the boundary and we characterize their limits in terms of the Green's function and of its regular part. Moreover, when the domain is a ball we prove that the limit concentration points are antipodal with respect to the center of the ball, the solution is axially symmetric with respect to the line joining these points and the nodal surface intersects the boundary. In the second case, i.e. when (1.9) holds, we are able to prove that, if  $n \geq 4$ , the positive and negative part of the solution concentrate at the same point (i.e. we have "bubble tower solutions") and we get a precise estimate of the blow up rates, in terms of the distance of the concentration

points from the boundary of the domain. Moreover, if  $n \geq 6$  we prove that the unique limit of the concentration points is away from the boundary and it is a critical point of the Robin's function. As far as we know this is the first time that the phenomenon of different concentration points converging to the same point is analyzed for critical exponent problems. Indeed this never happens in the case of positive solutions, see [14].

Note that solutions of both type exist, at least in symmetric domains. Indeed for the first type of solutions it is enough to take a positive solution in a symmetric cap of the domain and reflect it by antisymmetry. In the second case "bubble tower" solutions have been recently found by Pistoia and Weth [15].

To describe more precisely our results, we introduce some notations. We denote by  $G$  the Green's function of the Laplace operator defined by  $\forall x \in \Omega$

$$-\Delta G(x, \cdot) = c_n \delta_x \quad \text{in } \Omega, \quad G(x, \cdot) = 0 \quad \text{on } \partial\Omega,$$

where  $\delta_x$  is the Dirac mass at  $x$  and  $c_n = (n-2)\omega_n$ , with  $\omega_n$  denoting the area of the unit sphere of  $\mathbb{R}^n$ . We denote by  $H$  the regular part of  $G$ , that is,

$$H(x_1, x_2) = |x_1 - x_2|^{2-n} - G(x_1, x_2) \quad \text{for } (x_1, x_2) \in \Omega^2.$$

For  $x = (x_1, x_2) \in \Omega^2 \setminus \Gamma$ , with  $\Gamma = \{(y, y)/y \in \Omega\}$ , we denote by  $M(x)$  the matrix defined by

$$M(x) = (m_{ij})_{1 \leq i, j \leq 2}, \quad \text{where } m_{ii} = H(x_i, x_i), \quad m_{12} = m_{21} = G(x_1, x_2). \quad (1.4)$$

Then we have

**Theorem 1.2** *Let  $n \geq 3$  and let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1) and let  $a_{\varepsilon,1}, a_{\varepsilon,2}$  be the concentration points defined in Theorem 1.1. Assume that there exists a positive constant  $\eta$  such that*

$$\eta \leq -\max u_\varepsilon / \min u_\varepsilon \leq \eta^{-1}. \quad (1.5)$$

*Then,  $a_{\varepsilon,1}$  and  $a_{\varepsilon,2}$  are two global extremum points of  $u_\varepsilon$  and we have*

$$\mu_{\varepsilon,i} = |u_\varepsilon(a_{\varepsilon,i})|^{2/(n-2)-\varepsilon/2} / \beta_n^{2/(n-2)} \quad \text{for } i = 1, 2. \quad (1.6)$$

*In addition, there exists a positive constant  $\gamma$  such that, for  $\varepsilon$  small,*

$$|a_{\varepsilon,1} - a_{\varepsilon,2}| \geq \gamma, \quad d(a_{\varepsilon,i}, \partial\Omega) \geq \gamma \quad \text{for } i = 1, 2. \quad (1.7)$$

*More precisely, we have*

$$(c_2\varepsilon/c_1)^{1/2} |u_\varepsilon(a_{\varepsilon,i})| \rightarrow \beta_n / \bar{\Lambda}_i, \quad a_{\varepsilon,i} \rightarrow \bar{a}_i \in \Omega \quad \text{with } \bar{a}_1 \neq \bar{a}_2, \quad \text{for } i = 1, 2, \quad (1.8)$$

*where  $\bar{\Lambda}_i$  is a positive constant,*

$$c_1 = \beta_n \int_{\mathbb{R}^n} \delta_{(0,1)}^{(n+2)/(n-2)} \quad \text{and} \quad c_2 = \frac{n-2}{2} \beta_n^{2n/(n-2)} \int_{\mathbb{R}^n} \text{Log}(1 + |x|^2) \frac{(|x|^2 - 1)dx}{(1 + |x|^2)^{n+1}}.$$

*Furthermore,  $(\bar{a}_1, \bar{a}_2, \bar{\Lambda}_1, \bar{\Lambda}_2)$  is a critical point of the function*

$$\Psi : \Omega^2 \setminus \Gamma \times (0, \infty)^2 \rightarrow \mathbb{R}; \quad (a, \Lambda) := (a_1, a_2, \Lambda_1, \Lambda_2) \mapsto \frac{1}{2} {}^t \Lambda M(a) \Lambda - \log(\Lambda_1 \Lambda_2),$$

*where  $M(a)$  is the matrix defined by (1.4).*

**Remark 1.3** 1- *The assumption (1.5) allows us to prove that the distance between the concentration points is bounded from below by a positive constant.*

2- *The low energy positive solutions of (1) have to blow up at a critical point of the Robin's function  $\varphi(x) = H(x, x)$ , see [13] and [18].*

3- *A similar matrix to that involved in the above function  $\Psi$  plays also a crucial role in the characterization of the concentration points for the positive solutions of (1), see [3].*

The next result describes the asymptotic behavior, as  $\varepsilon \rightarrow 0$ , of low energy sign-changing solutions of (1) satisfying (1.5) outside the limit concentration points.

**Theorem 1.4** *Let  $n \geq 3$  and let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1) and (1.5). Then the limit concentration points  $\bar{a}_1$  and  $\bar{a}_2$ , defined in Theorem 1.2, are two isolated simple blow-up points of  $(u_\varepsilon)$  (see [14] for definition) and there exist positive constants  $m_1$  and  $m_2$  such that*

$$\varepsilon^{-1/2}u_\varepsilon \rightarrow \bar{u} := m_1G(\bar{a}_1, \cdot) - m_2G(\bar{a}_2, \cdot) \text{ in } C_{loc}^2(\bar{\Omega} \setminus \{\bar{a}_1, \bar{a}_2\}) \quad \text{as } \varepsilon \rightarrow 0.$$

The second part of this paper is devoted to the study of symmetry properties of low energy sign-changing solutions of (1) of the first type when  $\Omega$  is a ball. We shall prove the following results.

**Theorem 1.5** *Assume that  $n \geq 3$ . Let  $\Omega$  be the unit ball and let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1) and (1.5). Then, up to a rotation of  $\Omega$ , the limit concentration points  $\bar{a}_i$ 's and the reals  $\bar{\Lambda}_i$ 's, defined in Theorem 1.2, satisfy*

$$\begin{aligned} \bar{a}_1 &= -\bar{a}_2 = a_* := (0, \dots, 0, t_*), \\ \bar{\Lambda}_1 &= \bar{\Lambda}_2 = \left( \frac{1}{H(a_*, a_*) + G(a_*, -a_*)} \right)^{1/2}, \end{aligned}$$

where  $t_*$  is the unique solution of

$$g(t) = \frac{t}{(1-t^2)^{n-1}} - \frac{1}{(2t)^{n-1}} + \frac{t}{(1+t^2)^{n-1}} = 0, \quad \text{for } t \in (0, 1).$$

The characterization of the points  $\bar{a}_i$ 's and the reals  $\bar{\Lambda}_i$ 's allows us to improve the result of Theorem 1.4 and therefore we can prove that the nodal surface intersects the boundary. In fact, we have

**Theorem 1.6** *Under the assumptions of Theorem 1.5, we have*

- (a) *The constants  $m_1$  and  $m_2$  defined in Theorem 1.4 are equal.*
- (b) *The nodal surface of  $u_\varepsilon$  intersects the boundary  $\partial\Omega$ .*

Observe that the limit function  $G(\bar{a}_1, \cdot) - G(\bar{a}_2, \cdot)$ , defined in Theorem 1.4 with  $m_1 = m_2$ , is symmetric with respect to any hyperplane passing through the points  $\bar{a}_1$  and  $\bar{a}_2$ . Furthermore, it is antisymmetric with respect to the hyperplane passing through the origin and which is orthogonal to the line passing through the points  $\bar{a}_1$  and  $\bar{a}_2$ . Moreover, it changes sign once.

Following the idea of [10], we can prove, for  $\varepsilon$  small, that the functions  $u_\varepsilon$  satisfy also the symmetry property. More precisely, we have

**Theorem 1.7** *Let  $n \geq 3$  and let  $\Omega$  be a ball and  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1) and (1.5). Then, for  $\varepsilon$  sufficiently small, the concentration points  $a_{\varepsilon,1}$  and  $a_{\varepsilon,2}$  of  $u_\varepsilon$ , given by Theorem 1.1, are far away from the origin and they lay on the same line passing through the origin and  $u_\varepsilon$  is axially symmetry with respect to this line.*

*Moreover the points  $a_{\varepsilon,1}$  and  $a_{\varepsilon,2}$  lay on different sides with respect to  $T$  and all the critical points of  $u_\varepsilon$  belong to the symmetry axis and*

$$\frac{\partial u_\varepsilon}{\partial \nu_T}(x) > 0 \quad \forall x \in T \cap \Omega,$$

*where  $T$  is any hyperplane passing through the origin but not containing  $a_{\varepsilon,1}$  and where  $\nu_T$  is the normal to  $T$ , oriented towards the half space containing  $a_{\varepsilon,1}$ .*

Concerning the antisymmetric property, it is easy to construct a family of solutions  $(u_\varepsilon)$  which are antisymmetric by reflecting the positive minimizing solution on the half ball and hence interesting questions arise: Let  $(u_\varepsilon)$  be a family of solutions satisfying the assumptions of our Theorems, are the solutions  $(u_\varepsilon)$  antisymmetric? What about the uniqueness? A further investigation in this direction is in progress.

The last part of this paper is devoted to the study of the case where the assumption (1.5) is removed, that means the case when we can have sign-changing bubble tower solutions. We recall that Pistoia and Weth [15] have constructed such solutions in symmetric domains. Without loss of generality, we can assume, in the case where the assumption (1.5) is removed, that the following holds:

$$\frac{\max u_\varepsilon}{\min u_\varepsilon} \rightarrow -\infty \quad \text{as } \varepsilon \rightarrow 0. \quad (1.9)$$

Our main result reads:

**Theorem 1.8** *Let  $n \geq 4$  and let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1) and (1.9). Let  $a_{\varepsilon,1}$ ,  $a_{\varepsilon,2}$  be the concentration points and  $\mu_{\varepsilon,1}$ ,  $\mu_{\varepsilon,2}$  the speeds of the concentration points defined in Theorem 1.1. Then,*

*(a) there exists a positive constant  $c$  such that, for  $\varepsilon$  small,*

$$\begin{aligned} \frac{1}{c}\mu_{\varepsilon,1} \leq \mu_{\varepsilon,2}(\mu_{\varepsilon,2}d_{\varepsilon,2})^2 \leq c\mu_{\varepsilon,1}, \quad \frac{1}{c}\varepsilon \leq \frac{1}{(\mu_{\varepsilon,2}d_{\varepsilon,2})^{n-2}} \leq c\varepsilon, \quad \mu_{\varepsilon,1}\mu_{\varepsilon,2}|a_{\varepsilon,1} - a_{\varepsilon,2}|^2 \leq c, \\ \mu_{\varepsilon,2}|a_{\varepsilon,1} - a_{\varepsilon,2}| \rightarrow 0, \quad \frac{|a_{\varepsilon,1} - a_{\varepsilon,2}|}{d_{\varepsilon,2}} \rightarrow 0, \quad \frac{d_{\varepsilon,1}}{d_{\varepsilon,2}} \rightarrow 1 \quad \text{as } \varepsilon \rightarrow 0, \end{aligned}$$

*where  $d_{\varepsilon,i} = d(a_{\varepsilon,i}, \partial\Omega)$  for  $i = 1, 2$ .*

*(b) the nodal surface of  $u_\varepsilon$  does not intersect the boundary  $\partial\Omega$ .*

*Furthermore, if  $n \geq 6$ , we have*

$$\mu_{\varepsilon,1}\mu_{\varepsilon,2}|a_{\varepsilon,1} - a_{\varepsilon,2}|^2 \rightarrow 0, \quad d_{\varepsilon,i} \not\rightarrow 0 \quad \text{and} \quad a_{\varepsilon,i} \rightarrow \bar{a} \in \Omega \quad \text{for } i = 1, 2, \quad (1.10)$$

*where  $\bar{a}$  is a critical point of the Robin's function  $\varphi(x) = H(x, x)$ .*

Let us mention that we are not able to extend the results of Theorem 1.8 to the dimension 3 because of serious technical difficulties. Also the restriction to  $n \geq 6$  to deduce that the unique

limit concentration point is away from the boundary is due to some technical difficulties but we think that the same result should be true also in lower dimensions, at least for  $n = 4, 5$ . Moreover in the case of the ball since the only critical point of the Robin's function is the center, from Theorem 1.8 we deduce, for  $n \geq 6$ , that the bubble tower concentration point is the center, with negative and positive part blowing up with a prescribed rate. This makes us to conjecture that this solution should be radial and hence should be the only solution of this type.

The outline of the paper is the following. Section 2 is devoted to the proof of Theorems 1.1 and 1.2. We prove Theorems 1.4, 1.5, 1.6 and 1.7 in Section 3. Finally, Section 4 is devoted to the proof of Theorem 1.8.

## 2 Proof of Theorems 1.1 and 1.2

First, we deal with the proof of Theorem 1.1. Regarding the connected components of  $\Omega \setminus \{x \in \Omega \mid u_\varepsilon(x) = 0\}$ , let  $\Omega_1$  be one of them. Multiplying (1) by  $u_\varepsilon$  and integrating on  $\Omega_1$ , we derive that

$$\int_{\Omega_1} |\nabla u_\varepsilon|^2 \geq S^{n/2}(1 + o(1)), \quad (2.1)$$

where we have used Holder's inequality and the Sobolev embedding.

Since  $\|u_\varepsilon\|^2 \rightarrow 2S^{n/2}$  as  $\varepsilon$  goes to 0, we deduce that there are only two connected components.

The following lemma shows that the energy of the solution  $u_\varepsilon$  converges to  $S^{n/2}$  in each connected component. In fact we have

**Lemma 2.1** *Let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1). Then*

- (i)  $\int_{\Omega} |\nabla u_\varepsilon^+|^2 \rightarrow S^{n/2}$ ,  $\int_{\Omega} |\nabla u_\varepsilon^-|^2 \rightarrow S^{n/2}$  as  $\varepsilon \rightarrow 0$ ,
- (ii)  $\int_{\Omega} (u_\varepsilon^+)^{\frac{2n}{n-2}} \rightarrow S^{n/2}$ ,  $\int_{\Omega} (u_\varepsilon^-)^{\frac{2n}{n-2}} \rightarrow S^{n/2}$  as  $\varepsilon \rightarrow 0$ ,
- (iii)  $u_\varepsilon \rightarrow 0$  as  $\varepsilon \rightarrow 0$ ,
- (iv)  $M_{\varepsilon,+} := \max_{\Omega} u_\varepsilon^+ \rightarrow +\infty$ ,  $M_{\varepsilon,-} := \max_{\Omega} u_\varepsilon^- \rightarrow +\infty$  as  $\varepsilon \rightarrow 0$ ,

where  $u_\varepsilon^+ = \max(u_\varepsilon, 0)$  and  $u_\varepsilon^- = \max(0, -u_\varepsilon)$ .

**Proof.** The proof is the same as that of Lemma 2.1 of [8], so we omit it.  $\square$

Now, we are going to prove Theorem 1.1.

**Proof of Theorem 1.1** Arguing as in the proof of Theorem 1.1 of [8], we obtain that there exist  $a_{\varepsilon,1}$ ,  $a_{\varepsilon,2}$ ,  $\mu_{\varepsilon,1}$  and  $\mu_{\varepsilon,2}$  such that, as  $\varepsilon \rightarrow 0$ ,

$$\|u_\varepsilon^+ - P\delta_{(a_{\varepsilon,1}, \mu_{\varepsilon,1})}\| \rightarrow 0, \quad \|u_\varepsilon^- - P\delta_{(a_{\varepsilon,2}, \mu_{\varepsilon,2})}\| \rightarrow 0, \quad \mu_{\varepsilon,i} d(a_{\varepsilon,i}, \partial\Omega) \rightarrow +\infty, \text{ for } i = 1, 2. \quad (2.2)$$

Next, we will prove that one of the points is the global maximum point of  $|u_\varepsilon|$ . Arguing as in the proof of Lemma 2.2 of [8], we obtain

$$M_\varepsilon^{2/(n-2)-\varepsilon/2} d(a_\varepsilon, \partial\Omega) \rightarrow 0 \text{ as } \varepsilon \rightarrow 0, \quad (2.3)$$

where  $a_\varepsilon$  satisfies  $|u_\varepsilon(a_\varepsilon)| = |u_\varepsilon|_\infty := M_\varepsilon$ . Without loss of generality, we can assume that  $u_\varepsilon(a_\varepsilon) > 0$ . Now, let us define

$$\tilde{u}_\varepsilon(X) := M_\varepsilon^{-1} u_\varepsilon(a_\varepsilon + X/M_\varepsilon^{2/(n-2)-\varepsilon/2}), \text{ for } X \in \tilde{\Omega}_\varepsilon := M_\varepsilon^{2/(n-2)-\varepsilon/2}(\Omega - a_\varepsilon).$$

By Lemma 2.1, the limit domain of  $\tilde{\Omega}_\varepsilon$ , denoted by  $\Pi$ , has to be the whole space  $\mathbb{R}^n$  or a half space and by (2.3), it contains the origin. Since  $\tilde{u}_\varepsilon$  is bounded in  $\tilde{\Omega}_\varepsilon$ , using the standard elliptic theory, it converges in  $C_{loc}^2(\Pi)$  to a function  $u$  satisfying

$$-\Delta u = |u|^{2^*-2}u \text{ in } \Pi, \quad u(0) = 1, \quad u = 0 \text{ on } \partial\Pi, \quad \text{and} \quad \int_{\Pi} |\nabla u|^2 \leq 2S^{n/2}. \quad (2.4)$$

Observe that, any sign-changing solution  $w$  of (2.4) satisfies  $\|w\|^2 > 2S^{n/2}$ . Thus we derive that  $u$  is positive. It follows that  $\Pi$  has to be  $\mathbb{R}^n$  and  $u = \delta_{(0, \beta_n^{2/(2-n)})}$ . We deduce that  $\tilde{u}_\varepsilon > 0$  in any compact subset of  $\tilde{\Omega}_\varepsilon$ . But we have  $u_\varepsilon = 0$  on  $\partial\Omega_+$ , where  $\partial\Omega_+ := \{x \in \Omega : u_\varepsilon(x) > 0\}$ . Hence

$$\|u_\varepsilon^+ - M_\varepsilon^{\varepsilon(n-2)/2} P\delta_{(a_\varepsilon, \mu_\varepsilon)}\| \rightarrow 0 \text{ and } M_\varepsilon^{2/(n-2)-\varepsilon/2} d(a_\varepsilon, \partial\Omega_+) \rightarrow \infty \text{ as } \varepsilon \rightarrow 0, \quad (2.5)$$

where  $\mu_\varepsilon := M_\varepsilon^{2/(n-2)-\varepsilon/2} / \beta_n^{2/(n-2)}$ . Now, we observe that

$$\begin{aligned} M_\varepsilon^{\varepsilon(n-2)/2} S^{n/4} (1 + o(1)) &= M_\varepsilon^{\varepsilon(n-2)/2} \|P\delta_{(a_\varepsilon, \mu_\varepsilon)}\| \\ &\leq \|u_\varepsilon^+ - M_\varepsilon^{\varepsilon(n-2)/2} P\delta_{(a_\varepsilon, \mu_\varepsilon)}\| + \|u_\varepsilon^+\| \\ &= S^{n/4} (1 + o(1)). \end{aligned} \quad (2.6)$$

Thus, Lemma 2.1 and (2.6) imply that  $M_\varepsilon^\varepsilon$  goes to 1 as  $\varepsilon \rightarrow 0$ .

The proof of our theorem is thereby completed.  $\square$

The goal of the sequel of this section is to prove Theorem 1.2. We start by the following proposition which gives a parametrization of the function  $u_\varepsilon$ . It follows from the corresponding statements in [2].

**Proposition 2.2** *Let  $n \geq 3$  and let  $(u_\varepsilon)$  be a family of sign-changing solutions of (1) satisfying (1.1). Then the following minimization problem*

$$\min\{\|u_\varepsilon - \alpha_1 P\delta_{(a_1, \lambda_1)} + \alpha_2 P\delta_{(a_2, \lambda_2)}\|, \alpha_i > 0, \lambda_i > 0, a_i \in \Omega\}$$

has a unique solution  $(\alpha_1, \alpha_2, a_1, a_2, \lambda_1, \lambda_2)$  (up to permutation). In particular, we can write  $u_\varepsilon$  as follows

$$u_\varepsilon = \alpha_1 P\delta_{(a_1, \lambda_1)} - \alpha_2 P\delta_{(a_2, \lambda_2)} + v,$$

where  $v \in H_0^1(\Omega)$  such that  $\|v\| \rightarrow 0$  as  $\varepsilon \rightarrow 0$  and

$$(V_0) : \langle v, \varphi \rangle = 0 \text{ for } \varphi \in \{P\delta_{(a_i, \lambda_i)}, \partial P\delta_{(a_i, \lambda_i)}/\partial \lambda_i, \partial P\delta_{(a_i, \lambda_i)}/\partial a_i^j, i = 1, 2, 1 \leq j \leq n\}, \quad (2.7)$$

where  $a_i^j$  denotes the  $j$ -th component of  $a_i$ .



**Remark 2.3** For each  $i = 1, 2$ , the point  $a_i$  is close to  $a_{\varepsilon,i}$  and each parameter  $\lambda_i$  satisfies  $\lambda_i/\mu_{\varepsilon,i}$  is close to 1, where  $a_{\varepsilon,i}$ ,  $\mu_{\varepsilon,i}$  are defined in Theorem 1.1 and  $a_i$ ,  $\lambda_i$  are defined in Proposition 2.2.

As usual in this type of problems, we first deal with the  $v$ -part of  $u_\varepsilon$ , in order to show that it is negligible with respect to the concentration phenomena. Using the following estimate (see Subsection 3.1 of [19])

$$\delta_i^{-\varepsilon}(x) = \beta_n^{-\varepsilon} \lambda_i^{-\varepsilon(n-2)/2} + O(\varepsilon \log(1 + \lambda_i^2 |x - a_i|^2)) = 1 + o(1), \quad (2.8)$$

(since  $\lambda_i^\varepsilon \rightarrow 1$  as  $\varepsilon \rightarrow 0$ ) and arguing as in Lemma 3.3 of [8] we derive the following lemma.

**Lemma 2.4** The function  $v$  defined in Proposition 2.2 satisfies the following estimate

$$\|v\| \leq c\varepsilon + c \begin{cases} \sum_i \frac{1}{(\lambda_i d_i)^{n-2}} + \varepsilon_{12} (\log \varepsilon_{12}^{-1})^{(n-2)/n} & \text{if } n < 6 \\ \sum_i \frac{1}{(\lambda_i d_i)^{(n+2)/2 - \varepsilon(n-2)}} + \varepsilon_{12}^{(n+2)/2(n-2)} (\log \varepsilon_{12}^{-1})^{(n+2)/2n} & \text{if } n \geq 6, \end{cases}$$

where

$$\varepsilon_{12} = \left( \frac{\lambda_1}{\lambda_2} + \frac{\lambda_2}{\lambda_1} + \lambda_1 \lambda_2 |a_1 - a_2|^2 \right)^{(2-n)/2}.$$

Now, arguing as in the proof of Propositions 3.4, 3.5 and 3.6 of [8] and using (2.8) we obtain the following results.

**Proposition 2.5** Assume that  $n \geq 3$  and let  $\alpha_i$ ,  $a_i$  and  $\lambda_i$  be the variables defined in Proposition 2.2. We then have

$$1 - \frac{\alpha_i^{4/(n-2)-\varepsilon}}{\beta_n^\varepsilon \lambda_i^{\varepsilon(n-2)/2}} = O\left(\varepsilon + \frac{1}{(\lambda_i d_i)^{n-2}} + \varepsilon_{12} + \|v\|\right), \quad (2.9)$$

$$\begin{aligned} \alpha_i c_1 \frac{n-2}{2} \frac{H(a_i, a_i)}{\lambda_i^{n-2}} - \alpha_j c_1 \left( \lambda_i \frac{\partial \varepsilon_{12}}{\partial \lambda_i} + \frac{n-2}{2} \frac{H(a_1, a_2)}{(\lambda_1 \lambda_2)^{(n-2)/2}} \right) - \alpha_i \frac{n-2}{2} c_2 \varepsilon \\ \leq c\varepsilon^2 + c \begin{cases} \sum_{k=1,2} \frac{\log(\lambda_k d_k)}{(\lambda_k d_k)^n} + \varepsilon_{12}^{\frac{n}{n-2}} \log \varepsilon_{12}^{-1} & (\text{if } n \geq 4), \\ \sum_{k=1,2} \frac{1}{(\lambda_k d_k)^2} + \varepsilon_{12}^2 (\log \varepsilon_{12}^{-1})^{2/3} & (\text{if } n = 3), \end{cases} \end{aligned} \quad (2.10)$$

where  $i, j \in \{1, 2\}$  with  $i \neq j$ , and  $c_1, c_2$  are defined in Theorem 1.2.

**Proposition 2.6** Let  $\alpha_i$ ,  $a_i$  and  $\lambda_i$  be the variables defined in Proposition 2.2.

(a) For  $n \geq 4$ , we have

$$\begin{aligned} \alpha_i \frac{1}{\lambda_i^{n-1}} \frac{\partial H(a_i, a_i)}{\partial a_i} + 2 \frac{\alpha_j}{\lambda_i} \left( \frac{\partial \varepsilon_{12}}{\partial a_i} - \frac{\partial H}{\partial a_i}(a_1, a_2) \frac{1}{(\lambda_1 \lambda_2)^{(n-2)/2}} \right) \\ = O\left( \sum_{k=1,2} \frac{1}{(\lambda_k d_k)^n} + \varepsilon_{12}^{\frac{n}{n-2}} \log \varepsilon_{12}^{-1} + \varepsilon_{12} (\log \varepsilon_{12}^{-1})^{\frac{n-2}{n}} + \frac{\varepsilon}{(\lambda_i d_i)^{n-1}} \right), \end{aligned} \quad (2.11)$$

where  $i, j \in \{1, 2\}$  with  $i \neq j$ .

(b) For  $n = 3$ , we assume that  $a_1 \neq a_2$  then we have for  $i \in \{1, 2\}$

$$\frac{1}{\lambda_i^2} \frac{\partial H(a_i, a_i)}{\partial a_i} + 2 \frac{1}{\lambda_i} \left( \frac{\partial \varepsilon_{12}}{\partial a_i} - \frac{\partial H}{\partial a_i}(a_1, a_2) \frac{1}{(\lambda_1 \lambda_2)^{1/2}} \right) = o \left( \sum_{k=1,2} \frac{1}{(\lambda_k \eta_k)^2} \right), \quad (2.12)$$

where  $\eta_1, \eta_2$  are two positive parameters chosen such that

$$B(a_1, \eta_1) \cap B(a_2, \eta_2) = \emptyset \quad \text{and} \quad B(a_i, \eta_i) \subset \Omega, \quad \text{for } i = 1, 2, .$$

Note that the proof of Propositions 2.5 and 2.6 are based on some integral estimates proved in [1] and [17].

To deal with dimension 3, we prove the following lemma.

**Lemma 2.7** *Assume that  $n = 3$  and assume further (1.5) holds. Then there exists a positive constant  $\bar{c}_0$  such that*

$$|a_1 - a_2| \geq \bar{c}_0 \max(d_1, d_2),$$

where the points  $a_i$ 's are defined in Proposition 2.2 and  $d_i = d(a_i, \partial\Omega)$ .

**Proof.** Arguing by contradiction, we assume that  $|a_1 - a_2| = o(\max(d_1, d_2))$ . This implies that  $d_1/d_2 \rightarrow 1$  as  $\varepsilon \rightarrow 0$ . Now, choosing  $\eta_1 = \eta_2 = |a_1 - a_2|/4$ , we see that

$$B(a_1, \eta_1) \cap B(a_2, \eta_2) = \emptyset \quad \text{and} \quad B(a_i, \eta_i) \subset \Omega, \quad \text{for } i = 1, 2,$$

and

$$\begin{aligned} \frac{1}{(\lambda_i \eta_i)^2} &\leq \frac{c}{\lambda_1 \lambda_2 |a_1 - a_2|^2} \leq c \varepsilon_{12}^2, \\ \frac{1}{\lambda_i^2} \left| \frac{\partial H(a_i, a_i)}{\partial a_i} \right| &\leq \frac{c}{(\lambda_i d_i)^2} = o \left( \frac{1}{\lambda_1 \lambda_2 |a_1 - a_2|^2} \right) = o(\varepsilon_{12}^2), \\ \frac{1}{\lambda_i} \frac{1}{(\lambda_1 \lambda_2)^{1/2}} \left| \frac{\partial H(a_1, a_2)}{\partial a_i} \right| &\leq \frac{c}{(\lambda_i d_i)^2} = o(\varepsilon_{12}^2), \\ \frac{1}{\lambda_i} \frac{1}{(\lambda_1 \lambda_2)^{1/2}} \left| \frac{\partial}{\partial a_i} \left( \frac{1}{|a_1 - a_2|} \right) \right| &= \frac{1}{\lambda_i (\lambda_1 \lambda_2)^{1/2} |a_1 - a_2|^2} \geq c \varepsilon_{12}^2, \end{aligned}$$

where we have used the fact that  $\lambda_1$  and  $\lambda_2$  are of the same order. Applying (2.12) and the above estimates, we derive a contradiction and therefore our lemma follows.  $\square$

Next we prove the following crucial lemmas.

**Lemma 2.8** *Under the assumptions of Theorem 1.2, there exists a positive constant  $c_0 > 0$  such that the variable  $a_i$ , defined in Proposition 2.2, satisfy*

$$(i) \quad c_0 \leq \frac{d_1}{d_2} \leq c_0^{-1}; \quad (ii) \quad c_0 \leq \frac{|a_1 - a_2|}{d_i} \leq c_0^{-1}, \quad \text{for } i = 1, 2,$$

where  $d_i = d(a_i, \partial\Omega)$ .

**Proof.** On one hand, using (2.10), we have

$$\varepsilon = O\left(\sum \frac{1}{(\lambda_j d_j)^{n-2}} + \varepsilon_{12}\right). \quad (2.13)$$

On the other hand, using (1.5) an easy computation shows that

$$\varepsilon_{12} = \frac{1}{(\lambda_1 \lambda_2 |a_1 - a_2|^2)^{(n-2)/2}} + O(\varepsilon_{12}^{n/(n-2)}), \quad (2.14)$$

$$\lambda_i \frac{\partial \varepsilon_{12}}{\partial \lambda_i} := -\frac{n-2}{2} \varepsilon_{12} \left(1 - 2 \frac{\lambda_j}{\lambda_i} \varepsilon_{12}^{\frac{2}{n-2}}\right) = -\frac{n-2}{2} \frac{1}{(\lambda_1 \lambda_2 |a_1 - a_2|^2)^{(n-2)/2}} + o(\varepsilon_{12}). \quad (2.15)$$

Thus, using (2.9), (2.13) and (2.15), the estimate (2.10) becomes

$$\frac{H(a_i, a_i)}{\lambda_i^{n-2}} + \frac{G(a_1, a_2)}{(\lambda_1 \lambda_2)^{(n-2)/2}} - \frac{c_2}{c_1} \varepsilon = o\left(\varepsilon_{12} + \sum \frac{1}{(\lambda_j d_j)^{n-2}}\right). \quad (2.16)$$

Now we claim that

$$\frac{1}{\lambda_i^{n-1}} \frac{\partial H(a_i, a_i)}{\partial a_i} + \frac{2}{\lambda_i} \frac{\partial G}{\partial a_i}(a_1, a_2) \frac{1}{(\lambda_1 \lambda_2)^{(n-2)/2}} = o\left(\varepsilon_{12}^{(n-1)/(n-2)} + \sum \frac{1}{(\lambda_j d_j)^{n-1}}\right). \quad (2.17)$$

For  $n \geq 4$ , (2.17) follows immediately from (2.9), (2.11) and (2.13).

For  $n = 3$ , choosing  $\eta_i = \min(\bar{c}_0, 1)d_i/4$  in (2.12) where  $\bar{c}_0$  is the positive constant defined in Lemma 2.7, Claim (2.17) follows from (2.12).

Now we are going to prove Claim (i). Arguing by contradiction, we assume, for example, that  $d_2 = o(d_1)$ . Using (2.16) for  $i = 1$  and  $i = 2$ , we get

$$\frac{1}{(\lambda_2 d_2)^{n-2}} = o(\varepsilon_{12}). \quad (2.18)$$

Using (2.14), (2.18) and the fact that  $|\partial H(a_1, a_2)/\partial a_k| \leq cd_k^{-1}(d_1 d_2)^{(2-n)/2}$ , it is easy to obtain

$$\frac{1}{(\lambda_1 \lambda_2)^{(n-2)/2}} \left| \frac{1}{\lambda_k} \frac{\partial G}{\partial a_k}(a_1, a_2) \right| \geq c\varepsilon_{12}^{(n-1)/(n-2)} \quad \text{for } k = 1, 2. \quad (2.19)$$

Clearly, (2.17), (2.18) and (2.19) give a contradiction. Thus, we derive that  $d_1$  and  $d_2$  are of the same order. Hence Claim (i) is proved. Regarding Claim (ii), arguing by contradiction, we assume that  $d_1 = o(|a_1 - a_2|)$ . In this case, it is easy to obtain

$$\begin{aligned} \frac{1}{(\lambda_1 \lambda_2)^{(n-2)/2}} \left| \frac{\partial G}{\partial a_1}(a_1, a_2) \right| &\leq \frac{1}{(\lambda_1 \lambda_2)^{(n-2)/2}} \left( \frac{c}{|a_1 - a_2|^{n-1}} + \frac{c}{d_1 |a_1 - a_2|^{n-2}} \right) \\ &= o\left(\frac{1}{d_1 (\lambda_1 d_1)^{n-2}}\right). \end{aligned} \quad (2.20)$$

Thus, (2.17) and (2.20) give again a contradiction and we derive that  $d_1/|a_1 - a_2|$  is bounded below. Hence, by Lemma 2.7, the proof is completed for  $n = 3$ .

It remains to prove that  $d_1/|a_1 - a_2|$  is bounded above for  $n \geq 4$ . To this aim, we argue by contradiction and we assume that  $|a_1 - a_2| = o(d_1)$ . Therefore,

$$\begin{aligned} \frac{1}{\lambda_1^{n-1}} \left| \frac{\partial H(a_1, a_1)}{\partial a_1} \right| + \frac{1}{\lambda_1(\lambda_1 \lambda_2)^{(n-2)/2}} \left| \frac{\partial H(a_1, a_2)}{\partial a_1} \right| \\ \leq \frac{c}{(\lambda_1 d_1)^{n-1}} = o\left( \frac{1}{(\lambda_1 \lambda_2 |a_1 - a_2|^2)^{(n-1)/2}} \right). \end{aligned} \quad (2.21)$$

We now observe that

$$\left| \frac{1}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_1} \right| \geq \frac{c}{\lambda_1(\lambda_1 \lambda_2)^{(n-2)/2}} \frac{1}{|a_1 - a_2|^{n-1}} \geq \frac{c}{(\lambda_1 \lambda_2 |a_1 - a_2|^2)^{(n-1)/2}}. \quad (2.22)$$

Hence, (2.17), (2.21) and (2.22) give a contradiction, and therefore  $|a_1 - a_2|/d_1$  is bounded below. Finally, using Claim (i), the proof of Claim (ii) is completed.  $\square$

Now, we will prove that the concentration points are in a compact set of  $\Omega$  and they are far away of each other.

**Lemma 2.9** *There exists a positive constant  $d_0$  such that*

$$|a_1 - a_2| \geq d_0 \quad ; \quad d_i \geq d_0 \quad \text{for } i = 1, 2.$$

**Proof.** The proof is the same as that of Lemma 3.8 of [8], so we omit it.  $\square$

Now we are ready to prove Theorem 1.2.

**Proof of Theorem 1.2** Without loss of generality we can assume that

$$M_\varepsilon := \max u_\varepsilon \geq M_{\varepsilon,-} := -\min u_\varepsilon.$$

Hence (2.5) holds. Now, let  $b_\varepsilon$  be such that  $M_{\varepsilon,-} := -u_\varepsilon(b_\varepsilon)$ . Using (1.5) and arguing as in the proof of (2.5) we can prove that

$$\|u_\varepsilon^- - M_{\varepsilon,-}^{\varepsilon(n-2)/2} P\delta_{(b_\varepsilon, \mu_{\varepsilon,-})}\| \rightarrow 0 \quad \text{and} \quad M_{\varepsilon,-}^{2/(n-2)-\varepsilon/2} d(b_\varepsilon, \partial\Omega_-) \rightarrow \infty \quad \text{as } \varepsilon \rightarrow 0, \quad (2.23)$$

where  $\mu_{\varepsilon,-} := M_{\varepsilon,-}^{2/(n-2)-\varepsilon/2} / \beta_n^{2/(n-2)}$  and  $\Omega_- := \{x \in \Omega : u_\varepsilon(x) < 0\}$ . Hence  $a_{\varepsilon,1}$  and  $a_{\varepsilon,2}$  can be chosen as  $a_\varepsilon$  and  $b_\varepsilon$  which are two global extremum points of  $u_\varepsilon$ .

Regarding, Claim (1.7), it follows from Lemma 2.9. Therefore each  $a_i$  converges to  $\bar{a}_i \in \Omega$  with  $\bar{a}_1 \neq \bar{a}_2$ .

Now, let us introduce the following change of variable

$$\frac{1}{\lambda_i^{(n-2)/2}} = \Lambda_i \left( \frac{c_2 \varepsilon}{c_1} \right)^{1/2}.$$

Note that, (2.16) and (2.17) imply, for  $i, j = 1, 2$  with  $j \neq i$ ,

$$H(a_i, a_i) \Lambda_i + G(a_1, a_2) \Lambda_j - \frac{1}{\Lambda_i} = o(\Lambda_i), \quad (2.24)$$

$$\frac{\partial H(a_i, a_i)}{\partial a_i} \Lambda_i^2 + 2 \frac{\partial G(a_1, a_2)}{\partial a_i} \Lambda_1 \Lambda_2 = o(\Lambda_i^2). \quad (2.25)$$

Since each  $a_i$  converges to  $\bar{a}_i \in \Omega$  with  $\bar{a}_1 \neq \bar{a}_2$ , thus the functions  $H$ ,  $G$  and its derivatives are bounded. Therefore, from (2.24) and (1.5), it is easy to see that for each  $i = 1, 2$ ,  $\Lambda_i$  is bounded above and below. Hence, each  $\Lambda_i$  converges to  $\bar{\Lambda}_i > 0$  (up to a sequence) which implies (1.8) (see (1.6) and Remark 2.3). Passing to the limit in (2.24) and (2.25), we get

$$H(\bar{a}_i, \bar{a}_i)\bar{\Lambda}_i + G(\bar{a}_1, \bar{a}_2)\bar{\Lambda}_j - \bar{\Lambda}_i^{-1} = 0, \quad (2.26)$$

$$\frac{\partial H(\bar{a}_i, \bar{a}_i)}{\partial \bar{a}_i} \bar{\Lambda}_i^2 + 2 \frac{\partial G(\bar{a}_1, \bar{a}_2)}{\partial \bar{a}_i} \bar{\Lambda}_1 \bar{\Lambda}_2 = 0, \quad (2.27)$$

where  $i, j = 1, 2$  with  $j \neq i$ .

Equations (2.26) and (2.27) imply that  $\nabla \Psi(\bar{a}_1, \bar{a}_2, \bar{\Lambda}_1, \bar{\Lambda}_2) = 0$ . Hence  $(\bar{a}_1, \bar{a}_2, \bar{\Lambda}_1, \bar{\Lambda}_2)$  is a critical point of  $\Psi$ .  $\square$

### 3 Proof of Theorems 1.4, 1.5, 1.6 and 1.7

Regarding Theorem 1.5, it follows immediately from Theorem 1.2 and the following lemma.

**Lemma 3.1** *Let  $\Omega$  be a ball and assume that  $n \geq 3$ . Then, up to a rotation of  $\Omega$ , the function  $\Psi$ , defined in Theorem 1.2, has only one critical point  $\bar{X} := (a, b, x, y)$ . It satisfies*

$$\begin{aligned} a = -b &= (0, \dots, 0, t_*) \text{ with } t_* > 0, \\ x = y &= \left( \frac{1}{H(a, a) + G(a, -a)} \right)^{1/2}, \end{aligned}$$

where  $t_*$  is the unique solution of

$$g(t) = \frac{t}{(1-t^2)^{n-1}} - \frac{1}{(2t)^{n-1}} + \frac{t}{(1+t^2)^{n-1}} = 0, \quad \text{for } t \in (0, 1).$$

**Proof.** Let  $(a_1, a_2, x_1, x_2)$  be a critical point of  $\Psi$ . Then, for  $i, j = 1, 2$  with  $j \neq i$ , we derive

$$H(a_i, a_i)x_i + G(a_1, a_2)x_j = \frac{1}{x_i} \quad (3.1)$$

$$x_i \frac{\partial H(a, a)}{\partial a} \Big|_{a=a_i} + 2x_j \frac{\partial G(a, a_j)}{\partial a} \Big|_{a=a_i} = 0. \quad (3.2)$$

Multiplying (3.1) by  $x_i$ , we get

$$H(a_1, a_1)x_1^2 = H(a_2, a_2)x_2^2. \quad (3.3)$$

Recall that when  $\Omega$  is the unit ball, we have

$$G(a, b) = \frac{1}{|a-b|^{n-2}} - \frac{1}{(|a|^2|b|^2 + 1 - 2\langle a, b \rangle)^{(n-2)/2}} \quad (3.4)$$

$$H(a, a) = \frac{1}{(1-|a|^2)^{n-2}}. \quad (3.5)$$

Thus,

$$\frac{\partial H(a, a)}{\partial a} = \frac{2(n-2)a}{(1-|a|^2)^{n-1}} \quad (3.6)$$

$$\frac{\partial G}{\partial a}(a, b) = \frac{(n-2)(b-a)}{|a-b|^n} - \frac{(n-2)(b-|b|^2 a)}{(|a|^2|b|^2 + 1 - 2\langle a, b \rangle)^{n/2}}. \quad (3.7)$$

First, using (3.2), (3.6) and (3.7), it is easy to prove that  $a_i \neq 0$  for  $i = 1, 2$ .

Without loss of generality, we can assume that  $a_1 = (0, \dots, 0, \gamma_1)$ , where  $\gamma_1$  is a constant. Taken the  $j$ -th component (for  $j = 1, \dots, n-1$ ) of the vector defined by (3.2), with  $i = 1$ , it follows that  $a_2 = (0, \dots, 0, \gamma_2)$ , where  $\gamma_2$  is a constant. Hence  $a_1$  and  $a_2$  lay in the same line passing through the origin. It remains to prove that  $\gamma_1 = -\gamma_2$ .

Using (3.2), (3.3), (3.6) and (3.7), we get

$$\frac{\gamma_i}{(1-\gamma_i^2)} (H(a_1, a_1)H(a_2, a_2))^{1/2} + \frac{\gamma_j - \gamma_i}{|\gamma_2 - \gamma_1|^n} - \frac{\gamma_j}{(1-\gamma_2\gamma_1)^{n-1}} = 0, \quad (3.8)$$

for  $i, j = 1, 2$  with  $j \neq i$ .

Adding (3.8) for  $i = 1$  and  $i = 2$ , we derive

$$\frac{(\gamma_1 + \gamma_2)(1 - \gamma_1\gamma_2)}{(1 - \gamma_1^2)^{n/2}(1 - \gamma_2^2)^{n/2}} = \frac{\gamma_1 + \gamma_2}{(1 - \gamma_1\gamma_2)^{n-1}}. \quad (3.9)$$

Thus, if  $\gamma_1 + \gamma_2 \neq 0$ , (3.9) implies that  $(1 - \gamma_1\gamma_2)^2 = (1 - \gamma_1^2)(1 - \gamma_2^2)$  which implies that  $\gamma_1 = \gamma_2$  and therefore  $a_1 = a_2$  which is a contradiction. Thus  $\gamma_1 + \gamma_2 = 0$ , that means  $a_1 = -a_2 = (0, \dots, 0, t_*)$ , with  $t_*$  is the unique solution of

$$g(t) = \frac{t}{(1-t^2)^{n-1}} - \frac{1}{(2t)^{n-1}} + \frac{t}{(1+t^2)^{n-1}} = 0, \quad \text{for } t \in (0, 1),$$

where we have used (3.8).

Now using (3.3), (3.5) and the fact that the reals  $x_i$ 's are positive, it is easy to obtain that  $x_1 = x_2$ . Using again (3.1) we derive that

$$x_1 = x_2 = \left( \frac{1}{H(a_1, a_1) + G(a_1, a_2)} \right)^{1/2} \quad (3.10)$$

which completes the proof of our lemma.  $\square$

Next we are going to prove Theorem 1.4.

**Proof of Theorem 1.4** Observe that, by Theorems 1.1 and 1.2, we know that  $u_\varepsilon$  can be written as  $P\delta_{(a_{\varepsilon,1}, \mu_{\varepsilon,1})} - P\delta_{(a_{\varepsilon,2}, \mu_{\varepsilon,2})} + v$  with  $\|v\| \rightarrow 0$ ,  $u_\varepsilon(a_{\varepsilon,1}) = \max u_\varepsilon$ ,  $u_\varepsilon(a_{\varepsilon,2}) = \min u_\varepsilon$  and  $\mu_{\varepsilon,i}|a_{\varepsilon,1} - a_{\varepsilon,2}| \rightarrow \infty$ , for  $i = 1, 2$ . Furthermore, the concentration speeds satisfy (1.8).

Set  $h_\varepsilon := \max d(x, \mathcal{S})^{(n-2)/2} |u_\varepsilon(x)|$  where  $\mathcal{S} = \{a_{\varepsilon,1}, a_{\varepsilon,2}\}$ . It is easy to prove that  $h_\varepsilon$  is bounded (if not, we can construct another blow-up point and therefore the energy of  $u_\varepsilon$  becomes bigger than  $3S^{n/2}$  which gives a contradiction).

Let  $d_{\varepsilon,1} = d(a_{\varepsilon,1}, \partial\Omega_+)$  and  $d_{\varepsilon,2} = d(a_{\varepsilon,2}, \partial\Omega_-)$ , where  $\Omega_+ = \{x \in \Omega : u_\varepsilon(x) > 0\}$  and  $\Omega_- =$

$\{x \in \Omega : u_\varepsilon(x) < 0\}$ . We need to prove that  $d_{\varepsilon,i} \not\rightarrow 0$  as  $\varepsilon \rightarrow 0$ . Arguing by contradiction, assume that  $d_{\varepsilon,1} \leq d_{\varepsilon,2}$  and  $d_{\varepsilon,1} \rightarrow 0$ . We define the following function

$$w_\varepsilon(X) := d_{\varepsilon,1}^{\alpha_\varepsilon} u_\varepsilon(a_{\varepsilon,1} + d_{\varepsilon,1}X) \quad \text{for } X \in \Omega'_{\varepsilon,+} := d_{\varepsilon,1}^{-1}(\Omega_+ - a_{\varepsilon,1}),$$

where  $\alpha_\varepsilon = 2(n-2)/(4-\varepsilon(n-2))$ . An easy computation shows that

$$-\Delta w_\varepsilon = w_\varepsilon^{p-\varepsilon}, \quad w_\varepsilon > 0 \quad \text{in } \Omega'_{\varepsilon,+}, \quad w_\varepsilon = 0 \quad \text{on } \partial\Omega'_{\varepsilon,+}.$$

Observe that  $B(0,1) \subset \Omega'_{\varepsilon,+}$  and  $w_\varepsilon > 0$  in  $\Omega'_{\varepsilon,+}$ . Since  $h_\varepsilon$  is bounded and  $d_{\varepsilon,1}^{(n-2)/2} u_\varepsilon(a_{\varepsilon,1}) \rightarrow \infty$  (see (2.5) and (2.23)), we derive that 0 is an isolated blow-up point of  $(w_\varepsilon)$ . Thus, using [14], we deduce that 0 is isolated simple blow-up point of  $(w_\varepsilon)$ . Hence, we have

$$w_\varepsilon(0)w_\varepsilon(y) \leq c|y|^{2-n}, \quad \text{for all } |y| \leq 1/2. \quad (3.11)$$

By standard elliptic theories, we derive that  $w_\varepsilon(0)w_\varepsilon$  converges in  $C_{loc}^2(\Pi)$  to a function  $w$  satisfying

$$-\Delta w = 0 \quad \text{in } \Pi \setminus \{0\}, \quad w = 0 \quad \text{on } \partial\Pi,$$

where  $\Pi$  is the limit domain of  $\Omega'_{\varepsilon,+}$ . Since 0 is an isolated simple blow-up point of  $(w_\varepsilon)$  we deduce that 0 is a nonremovable singularity and therefore  $w = cG_\Pi$ , where  $G_\Pi$  is the Green's function and  $c$  is a positive constant. Now, using Pohozaev identity in the form of Corollary 1.1 of [14] we obtain

$$\underline{c}\varepsilon(1+o(1)) \int_{B(0,\sigma)} w_\varepsilon^{p+1-\varepsilon} - \frac{\sigma}{p+1-\varepsilon} \int_{\partial B(0,\sigma)} w_\varepsilon^{p+1-\varepsilon} = \int_{\partial B(0,\sigma)} B(\sigma, x, w_\varepsilon, \nabla w_\varepsilon) dx, \quad (3.12)$$

where  $\underline{c}$  is a positive constant and

$$B(\sigma, x, w_\varepsilon, \nabla w_\varepsilon) = \frac{n-2}{2} w_\varepsilon \frac{\partial w_\varepsilon}{\partial \nu} - \frac{\sigma}{2} |\nabla w_\varepsilon|^2 + \sigma \left( \frac{\partial w_\varepsilon}{\partial \nu} \right)^2.$$

Observe that, using (3.11), we obtain

$$w_\varepsilon^2(0)\sigma \int_{\partial B(0,\sigma)} w_\varepsilon^{p+1-\varepsilon} \leq c w_\varepsilon(0)^{1+\varepsilon-p} \sigma^{\varepsilon(n-2)-n} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0, \quad (3.13)$$

$$\varepsilon(1+o(1))w_\varepsilon^2(0) \int_{B(0,\sigma)} w_\varepsilon^{p+1-\varepsilon} \sim c d_{\varepsilon,1}^{2\alpha_\varepsilon} \varepsilon u_\varepsilon(a_{\varepsilon,1})^2 \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0, \quad (3.14)$$

where we have used (1.8) and the fact that  $\alpha_\varepsilon \rightarrow (n-2)/2$  and  $d_{\varepsilon,1} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ .

For the last term in (3.12), an easy computation shows

$$\lim_{\varepsilon \rightarrow 0, \sigma \rightarrow 0} \int_{\partial B(0,\sigma)} B(\sigma, x, w_\varepsilon(0)w_\varepsilon, w_\varepsilon(0)\nabla w_\varepsilon) dx = cH_\Pi(0,0). \quad (3.15)$$

Clearly, (3.12),..., (3.15) and the fact that  $\Pi \neq \mathbb{R}^n$  yield a contradiction. Hence  $d_{\varepsilon,1} \not\rightarrow 0$  as  $\varepsilon \rightarrow 0$  and therefore  $\bar{a}_1$  is an isolated simple blow up point of  $(u_\varepsilon)$ . The same holds for  $\bar{a}_2$ .

Now, arguing as in the proof of (4.10) of [7], the result follows.  $\square$

Now, we are going to prove Theorem 1.6.

**Proof of Theorem 1.6** We start by proving Claim (a). By Theorem 1.4, the points  $\bar{a}_i$ 's are two isolated simple blow-up points of  $(u_\varepsilon)$ . Thus, as in (3.12), we derive that

$$\underline{c}\varepsilon(1 + o(1)) \int_{B(a_{\varepsilon,i},\sigma)} u_\varepsilon^{p+1-\varepsilon} - \frac{\sigma}{p+1-\varepsilon} \int_{\partial B(a_{\varepsilon,i},\sigma)} u_\varepsilon^{p+1-\varepsilon} = \int_{\partial B(a_{\varepsilon,i},\sigma)} B(\sigma, x, u_\varepsilon, \nabla u_\varepsilon) dx, \quad (3.16)$$

for  $i = 1, 2$ , where  $\underline{c}$  is a positive constant independent of  $i$ . As in (3.13), we have

$$\sigma u_\varepsilon^2(a_{\varepsilon,i}) \int_{\partial B(a_{\varepsilon,i},\sigma)} u_\varepsilon^{p+1-\varepsilon} \leq c u_\varepsilon(a_{\varepsilon,i})^{1+\varepsilon-p} \sigma^{\varepsilon(n-2)-n} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0, \quad (3.17)$$

and using (1.8), we have

$$\varepsilon u_\varepsilon^2(a_{\varepsilon,i}) \int_{B(a_{\varepsilon,i},\sigma)} u_\varepsilon^{p+1-\varepsilon} \sim S^{n/2} \varepsilon u_\varepsilon(a_{\varepsilon,i})^2 \rightarrow S^{n/2} \frac{c_1 \beta_n^2}{c_2 \bar{\Lambda}_i^2} \quad \text{as } \varepsilon \rightarrow 0. \quad (3.18)$$

It remains to study the right side integral of (3.16). Using again (1.8) and Theorem 1.4, we derive that

$$u_\varepsilon(a_{\varepsilon,i}) u_\varepsilon \rightarrow m'_i G(\bar{a}_i, \cdot) - m'_j G(\bar{a}_j, \cdot) \quad (3.19)$$

where  $j \neq i$  and  $m'_1, m'_2$  are two positive constants satisfying  $m'_1/m'_2 = m_1/m_2$ . Thus

$$\lim_{\varepsilon \rightarrow 0, \sigma \rightarrow 0} \int_{\partial B(a_{\varepsilon,i},\sigma)} B(\sigma, x, u_\varepsilon(a_{\varepsilon,i}) u_\varepsilon, u_\varepsilon(a_{\varepsilon,i}) \nabla u_\varepsilon) dx = c'_1 (m'_i H(\bar{a}_i, \bar{a}_i) - m'_j G(\bar{a}_1, \bar{a}_2)), \quad (3.20)$$

for  $i, j = 1, 2$  with  $i \neq j$  and where  $c'_1$  is a positive constant independent of  $i$  and  $j$ .

Using (3.19) and (3.20) and the fact that  $\bar{\Lambda}_1 = \bar{\Lambda}_2$  (see Theorem 1.5), we deduce that

$$m'_1 H(\bar{a}_1, \bar{a}_1) - m'_2 G(\bar{a}_1, \bar{a}_2) = m'_2 H(\bar{a}_2, \bar{a}_2) - m'_1 G(\bar{a}_1, \bar{a}_2).$$

Hence since  $\bar{a}_1 = -\bar{a}_2$ , using (3.5), we derive that

$$(m'_1 - m'_2) (H(\bar{a}_1, \bar{a}_1) - G(\bar{a}_1, -\bar{a}_1)) = 0.$$

It is easy to verify that  $H(\bar{a}_1, \bar{a}_1) \neq G(\bar{a}_1, -\bar{a}_1)$  and therefore we obtain that  $m'_1 = m'_2$  which implies that  $m_1 = m_2$ . The proof of Claim (a) is thereby completed.

It remains to prove Claim (b). Arguing by contradiction and assuming that the set  $\{x \in \Omega, u_\varepsilon(x) = 0\}$  does not intersect the boundary  $\partial\Omega$ . Thus  $\partial u_\varepsilon / \partial \nu$  does not change sign which implies that  $\partial \bar{u} / \partial \nu$  does not change sign also, where  $\bar{u}$  is defined in Theorem 1.4. Now, since  $m_1 = m_2$  (see Claim (a)), an easy computation shows that

$$\int_{\partial\Omega} \frac{\partial \bar{u}}{\partial \nu}(x) dx = m_1 \int_{\partial\Omega} \frac{\partial}{\partial \nu} (G(\bar{a}_1, x) - G(\bar{a}_2, x)) dx = 0,$$

which implies a contradiction. Thus the result follows.  $\square$

Now, we are going to prove Theorem 1.7.

**Proof of Theorem 1.7** According to Theorem 1.5, we know that, for  $\varepsilon$  close to 0, both points



$a_{\varepsilon,i}$  are far away from the origin and they lay on different sides with respect to  $T$ , where  $T$  is any hyperplane passing through the origin but not containing  $a_{\varepsilon,1}$ . Arguing now as in the proof of Lemma 4.1 of [8], we see that the points  $a_{\varepsilon,i}$  lay on the same line passing through the origin. Lastly, the proof of the other statements of Theorem 1.7 is exactly the same as that of Theorem 1.5 of [8], so we omit it. This ends the proof of our result.  $\square$

## 4 Proof of Theorem 1.8

Throughout this section,  $c$  stands for a generic constant depending only on  $n$  and whose value may change in every step of the computations.

To prove Theorem 1.8, we need a delicate analysis and careful estimates. First, for  $\varepsilon$  sufficiently small, Proposition 2.2 implies that  $u_\varepsilon$  can be uniquely written as

$$u_\varepsilon = \alpha_{1,\varepsilon} P\delta_{(a_{1,\varepsilon}, \lambda_{1,\varepsilon})} - \alpha_{2,\varepsilon} P\delta_{(a_{2,\varepsilon}, \lambda_{2,\varepsilon})} + v_\varepsilon, \quad (4.1)$$

where  $v_\varepsilon \in H_0^1(\Omega)$ ,  $\|v_\varepsilon\| \rightarrow 0$  as  $\varepsilon \rightarrow 0$  and  $v_\varepsilon$  satisfies (2.7).

To simplify the notations, we write  $\alpha_i$ ,  $\lambda_i$ ,  $P\delta_i$  and  $v$  instead of  $\alpha_{i,\varepsilon}$ ,  $\lambda_{i,\varepsilon}$ ,  $P\delta_{(a_{i,\varepsilon}, \lambda_{i,\varepsilon})}$  and  $v_\varepsilon$  respectively. We will also use the following notations:

$$d_i = d(a_i, \partial\Omega) \quad \text{and} \quad \varepsilon_{12} = \left( \frac{\lambda_1}{\lambda_2} + \frac{\lambda_2}{\lambda_1} + \lambda_1 \lambda_2 |a_1 - a_2|^2 \right)^{(2-n)/2}.$$

Now, we prove the following crucial lemmas:

**Lemma 4.1** *Assume that  $n \geq 4$  and*

$$\varepsilon_{12} = (\lambda_1 \lambda_2 |a_1 - a_2|^2)^{(2-n)/2} + o(\varepsilon_{12}) \quad \text{as } \varepsilon \rightarrow 0. \quad (4.2)$$

*Then, for  $\varepsilon$  small, there exists a positive constant  $c$  such that*

$$(i) \quad \frac{d_1}{c} \leq d_2 \leq cd_1 \quad \text{and} \quad (ii) \quad \frac{\lambda_1}{c} \leq \lambda_2 \leq c\lambda_1.$$

**Proof.** On one hand, using (4.2), we derive that (2.15) holds. Thus, using (2.9), (2.10), (2.13) and (2.15), we see that (2.16) holds. On the other hand, (2.9), (2.11) and (2.13) imply that (2.17) holds. Now, arguing by contradiction we assume, for example, that  $\lambda_2 d_2 = o(\lambda_1 d_1)$ . Using (2.16) for  $i = 1$  and  $i = 2$ , we obtain (2.18). Using (4.2) and (2.18), we get

$$\frac{1}{(\lambda_1 \lambda_2)^{(n-2)/2}} \left| \frac{1}{\lambda_2} \frac{\partial G}{\partial a_2}(a_1, a_2) \right| \geq c\varepsilon_{12}^{(n-1)/(n-2)}. \quad (4.3)$$

Clearly, (2.16) for  $i = 2$ , (4.2), (2.18) and (4.3) give a contradiction. Thus, we derive that  $\lambda_1 d_1$  and  $\lambda_2 d_2$  are of the same order. Assume, for example, that  $d_1 = o(d_2)$ . In this case, it is easy to obtain that  $|a_1 - a_2| \geq d_2 - d_1 \geq d_2/2$ . Thus (2.20) holds. Obviously, (2.16) and (2.20) give a contradiction and therefore Claims (i) and (ii) are proved.  $\square$

**Lemma 4.2** *Let  $n \geq 4$  and assume that (1.9) holds. Then, for  $\varepsilon$  small, we have*

$$|a_1 - a_2| = o(d_2) \quad \text{and} \quad d_1 = d_2 + o(d_2).$$

**Proof.** Observe that Remark 2.3 implies that  $\lambda_1/\lambda_2 \rightarrow +\infty$  as  $\varepsilon \rightarrow 0$ . Thus, using Lemma 4.1, we derive that there exists a positive constant  $c$  such that

$$\frac{\lambda_1}{\lambda_2} \geq c\lambda_1\lambda_2 |a_1 - a_2|^2.$$

Thus  $\lambda_2 |a_1 - a_2| \leq c$ . Now, since  $\lambda_2 d_2 \rightarrow +\infty$  as  $\varepsilon \rightarrow 0$ , we derive that  $|a_1 - a_2| = o(d_2)$  and therefore  $d_1 = d_2 + o(d_2)$ .  $\square$

**Remark 4.3** *Notice that assumption  $\lambda_1/\lambda_2 \rightarrow +\infty$  as  $\varepsilon \rightarrow 0$  and Lemma 4.1 imply that*

$$\varepsilon_{12} \geq c \left( \frac{\lambda_2}{\lambda_1} \right)^{(n-2)/2} \quad \text{for } \varepsilon \text{ small.}$$

**Lemma 4.4** *Let  $n \geq 4$  and assume that (1.9) holds. Then, there exists a positive constant  $c$  such that, for  $\varepsilon$  small,*

$$\begin{aligned} (i) \quad & \frac{1}{c}\varepsilon_{12} \leq \frac{1}{(\lambda_2 d_2)^{n-2}} \leq c\varepsilon_{12} \\ (ii) \quad & \frac{1}{c}\varepsilon \leq \varepsilon_{12} \leq c\varepsilon \\ (iii) \quad & \frac{1}{c}\lambda_1 \leq \lambda_2(\lambda_2 d_2)^2 \leq c\lambda_1. \end{aligned}$$

**Proof.** Using (2.10) for  $i = 1$ , we see that

$$c_1\varepsilon_{12} - c_2\varepsilon = o\left(\frac{1}{(\lambda_2 d_2)^{n-2}} + \varepsilon_{12} + \varepsilon\right).$$

Thus

$$\varepsilon = \frac{c_1}{c_2}\varepsilon_{12} + o\left(\frac{1}{(\lambda_2 d_2)^{n-2}} + \varepsilon_{12}\right). \quad (4.4)$$

Using (2.13), (2.15) and (2.10) for  $i = 1$  and for  $i = 2$ , we derive that

$$c_1 \frac{H(a_2, a_2)}{\lambda_2^{n-2}} - 2c_1\varepsilon_{12}\varepsilon_{12}^{\frac{2}{n-2}} \frac{\lambda_1}{\lambda_2} (1 + o(1)) = o\left(\frac{H(a_2, a_2)}{\lambda_2^{n-2}} + \varepsilon_{12}\right). \quad (4.5)$$

But, by Remark 4.3, we know that there exists  $c > 0$  such that

$$c \leq \frac{\lambda_1}{\lambda_2} \varepsilon_{12}^{2/n-2} \leq 1. \quad (4.6)$$

Clearly, (4.4), (4.5), (4.6) and the fact that  $c \leq H(a_2, a_2)d_2^{n-2} \leq 1$  imply Claims (i) and (ii). Now, it follows from Claim (i) and Lemma 4.1 that the following holds:

$$c'' \left( \frac{\lambda_2}{\lambda_1} \right)^{\frac{n-2}{2}} \leq c'\varepsilon_{12} \leq \frac{1}{(\lambda_2 d_2)^{n-2}} \leq c\varepsilon_{12} \leq c \left( \frac{\lambda_2}{\lambda_1} \right)^{\frac{n-2}{2}},$$

where  $c$ ,  $c'$  and  $c''$  are positive constants.  
Therefore Claim (iii) follows.  $\square$

**Lemma 4.5** *Let  $n \geq 4$  and assume that (1.9) holds. Then, there exists a positive constant  $c$  such that, for  $\varepsilon$  small,*

$$\begin{aligned} (i) \quad & \lambda_1 \lambda_2 |a_1 - a_2|^2 \leq c \\ (ii) \quad & \lambda_2 |a_1 - a_2| \leq c \varepsilon_{12}^{1/n-2} \rightarrow 0 \quad \text{as } \varepsilon \rightarrow 0. \end{aligned}$$

**Proof.** Using (2.11) for  $i = 2$  and Lemma 4.4, we obtain

$$\frac{1}{\lambda_2^{n-1}} \frac{\partial H}{\partial a_2}(a_2, a_2) + 2(n-2)\lambda_1(a_1 - a_2)\varepsilon_{12}^{\frac{n}{n-2}} = o\left(\frac{1}{(\lambda_2 d_2)^{n-1}}\right) = o\left(\varepsilon_{12}^{\frac{n-1}{n-2}}\right). \quad (4.7)$$

Now, from (4.6) we deduce that

$$c\sqrt{\lambda_1 \lambda_2} |a_1 - a_2| \varepsilon_{12}^{\frac{n-1}{n-2}} \leq \lambda_1 |a_1 - a_2| \varepsilon_{12}^{n/(n-2)} \leq \sqrt{\lambda_1 \lambda_2} |a_1 - a_2| \varepsilon_{12}^{\frac{n-1}{n-2}}. \quad (4.8)$$

Arguing by contradiction, assume that  $\lambda_1 \lambda_2 |a_1 - a_2|^2 \rightarrow +\infty$ . Using (4.7), (4.8) and the fact that

$$\left| \frac{1}{\lambda_2^{n-1}} \frac{\partial H}{\partial a_2}(a_2, a_2) \right| \leq \frac{c}{(\lambda_2 d_2)^{n-1}} \leq c \varepsilon_{12}^{\frac{n-1}{n-2}},$$

we obtain a contradiction and therefore Claim (i) follows. Finally, Claim (ii) follows from Claim (i) and Remark 4.3.  $\square$

Now, to deal with the case of  $n \geq 6$ , we need the following crucial proposition which improves the estimate (2.11).

**Proposition 4.6** *Let  $n \geq 6$  and assume that (1.9) holds. Then we have*

$$\frac{1}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_1} = o\left(\varepsilon_{12}^{(n+1)/(n-2)}\right).$$

**Proof.** For sake of simplicity, we will use the following notations:

$$\begin{aligned} P\delta_i &= P\delta_{(a_i, \lambda_i)}, \quad \delta_i = \delta_{(a_i, \lambda_i)}, \quad \theta_i = \delta_i - P\delta_i, \\ \bar{u}_\varepsilon &= \alpha_1 P\delta_1 - \alpha_2 P\delta_2, \quad \varphi_1 = \frac{1}{\lambda_1} \frac{\partial P\delta_1}{\partial a_1}, \quad \psi_1 = \frac{1}{\lambda_1} \frac{\partial \delta_1}{\partial a_1}. \end{aligned}$$

Multiplying (1) by  $\varphi_1$  and integrating on  $\Omega$ , we obtain

$$\alpha_1 \int_{\Omega} \delta_1^p \varphi_1 - \alpha_2 \int_{\Omega} \delta_2^p \varphi_1 = \int_{\Omega} |u_\varepsilon|^{p-1-\varepsilon} u_\varepsilon \varphi_1, \quad (4.9)$$

where  $p = (n + 2)/(n - 2)$ .

For the left-hand side of (4.9), it follows from [1] that

$$\int_{\Omega} \delta_1^p \frac{1}{\lambda_1} \frac{\partial P \delta_1}{\partial a_1} = -\frac{1}{2} \frac{c_1}{\lambda_1^{n-1}} \frac{\partial H(a_1, a_1)}{\partial a_1} + O\left(\frac{1}{(\lambda_1 d_1)^n}\right), \quad (4.10)$$

$$\begin{aligned} \int_{\Omega} \delta_2^p \frac{1}{\lambda_1} \frac{\partial P \delta_1}{\partial a_1} &= \int_{\mathbb{R}^n} \delta_2^p \frac{1}{\lambda_1} \frac{\partial \delta_1}{\partial a_1} + O\left(\int_{\Omega^c} \delta_2^p \delta_1 + \int_{\Omega} \delta_2^p \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right|\right) \\ &= \frac{c}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_1} + O\left(\lambda_2 \left| a_1 - a_2 \right| \varepsilon_{12}^{\frac{n+1}{n-2}}\right) + o\left(\frac{1}{(\lambda_1 d_1)^{\frac{n-2}{2}} (\lambda_2 d_2)^{\frac{n}{2}}}\right) \\ &= \frac{c}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_1} + o\left(\varepsilon_{12}^{(n+1)/(n-2)}\right), \end{aligned} \quad (4.11)$$

where we have used Lemmas 4.2, 4.4 and 4.5. For the other integral in (4.9), an easy expansion implies that

$$\begin{aligned} \int_{\Omega} |u_{\varepsilon}|^{p-1-\varepsilon} u_{\varepsilon} \varphi_1 &= \int_{\Omega} |\bar{u}_{\varepsilon}|^{p-1-\varepsilon} \bar{u}_{\varepsilon} \varphi_1 + p \int_{\Omega} |\bar{u}_{\varepsilon}|^{p-1} v \varphi_1 \\ &\quad + O\left(\|v\|^2 + \varepsilon \|v\| + \|v\|^{p-\varepsilon} \left(\|\theta_1\| + \varepsilon_{12}^{1/2} (\log \varepsilon_{12}^{-1})^{\frac{n-2}{2n}}\right)\right) \\ &= \int_{\Omega} |\bar{u}_{\varepsilon}|^{p-1-\varepsilon} \bar{u}_{\varepsilon} \varphi_1 + p \int_{\Omega} |\bar{u}_{\varepsilon}|^{p-1} v \varphi_1 + o\left(\varepsilon_{12}^{(n+1)/(n-2)}\right), \end{aligned} \quad (4.12)$$

where we have used in the last inequality Lemmas 2.4 and 4.4 and the fact that  $\|\theta_1\| \leq c(\lambda_1 d_1)^{(2-n)/2}$ . For the last integral in (4.12), arguing as in (323) – (326) in [8], we obtain

$$\int_{\Omega} |\bar{u}_{\varepsilon}|^{p-1} v \varphi_1 \leq c \|v\| \left( \varepsilon_{12}^{\frac{n+2}{2(n-2)}} (\log \varepsilon_{12}^{-1})^{\frac{n+2}{2n}} + \frac{\log(\lambda_1 d_1)}{(\lambda_1 d_1)^{(n+2)/2}} \right) = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right), \quad (4.13)$$

where we have used in the last inequality Lemmas 2.4 and 4.4. It remains to study the first integral in the right-hand side of (4.12). Denoting  $f = \alpha_1 \delta_1 - \alpha_2 \delta_2$ , we observe that

$$\begin{aligned} \int_{\Omega} |\bar{u}_{\varepsilon}|^{p-1-\varepsilon} \bar{u}_{\varepsilon} \varphi_1 &= \int_{\mathbb{R}^n} |f|^{p-1-\varepsilon} f \psi_1 + (p - \varepsilon) \int_{\Omega} |f|^{p-1-\varepsilon} (\alpha_1 \theta_1 - \alpha_2 \theta_2) \varphi_1 \\ &\quad + O\left(\int_{\Omega} |f|^{p-\varepsilon} \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right| + \int_{2(\theta_1 + \theta_2) \leq |f|} |f|^{p-2} (\theta_1^2 + \theta_2^2) |\varphi_1|\right) \\ &\quad + O\left(\frac{1}{(\lambda_1 d_1)^{\frac{n-2}{2}} (\lambda_2 d_2)^{\frac{n+2}{2}}}\right) + O\left(\int (\theta_1^p + \theta_2^p) \left(|\psi_1| + \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right|\right)\right). \end{aligned} \quad (4.14)$$

Now, using Lemma 4.4, we can write

$$\int (\theta_1^p + \theta_2^p) \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right| \leq \left[ \frac{1}{(\lambda_1 d_1)^{\frac{n+2}{2}}} + \frac{1}{(\lambda_2 d_2)^{\frac{n+2}{2}}} \right] \frac{1}{(\lambda_1 d_1)^{\frac{n}{2}}} = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right); \quad (4.15)$$

$$\begin{aligned} \int \theta_1^p |\psi_1| &\leq \int_{B(a_1, d_1)} \theta_1^{n/(n-2)} \delta_1^{n/(n-2)} + \int_{B(a_1, d_1)^c} \delta_1^{p+1} \\ &\leq c \frac{\log(\lambda_1 d_1)}{(\lambda_1 d_1)^n} = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right); \end{aligned} \quad (4.16)$$

$$\begin{aligned} \int \theta_2^p |\psi_1| &\leq \int_{B(a_1, d_1)} \theta_2^{p/2} \delta_2^{p/2} |x - a_1| \delta_1^{n/(n-2)} + \int_{B(a_1, d_1)^c} \theta_2^p \delta_1 \\ &\leq \frac{c}{d_2^{\frac{n+2}{2}}} \frac{d_1^{\frac{n}{2}}}{\lambda_1^{\frac{n}{2}}} + \frac{c}{(\lambda_1 d_1)^{\frac{n-2}{2}} (\lambda_2 d_2)^{\frac{n+2}{2}}} = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right). \end{aligned} \quad (4.17)$$

Note that, since  $p - 2 \leq 0$ , we get

$$\int_{2(\theta_1 + \theta_2) \leq |f|} |f|^{p-2} (\theta_1^2 + \theta_2^2) |\varphi_1| \leq \int (\theta_1^p + \theta_2^p) (|\psi_1| + \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right|) = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right). \quad (4.18)$$

where we have used (4.15)-(4.17). Now, using Lemmas 4.2 and 4.4, we derive that

$$\begin{aligned} \int_{\Omega} |f|^{p-\varepsilon} \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right| + \int_{\Omega} |f|^{p-1-\varepsilon} \theta_1 |\varphi_1| &\leq \frac{1}{(\lambda_1 d_1^2)^{\frac{n-2}{2}}} \int (\delta_1^p + \delta_2^p) \\ &\leq \frac{c}{(\lambda_1 d_1)^{\frac{n-2}{2}} (\lambda_2 d_2)^{\frac{n-2}{2}}} = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right). \end{aligned} \quad (4.19)$$

Using again Lemmas 4.2 and 4.4, we obtain

$$\begin{aligned} \int_{\Omega} |f|^{p-1-\varepsilon} \theta_2 |\varphi_1| &\leq \int \delta_1^p \theta_2 + \int \delta_2^{p-1} \theta_2 |x - a_1| \delta_1^{\frac{n}{n-2}} + \int_{\Omega} \delta_2^{p-1} \theta_2 \frac{1}{\lambda_1} \left| \frac{\partial \theta_1}{\partial a_1} \right| \\ &\leq \frac{1}{(\lambda_1 d_1)^{\frac{n-2}{2}}} \frac{1}{(\lambda_2 d_2)^{\frac{n-2}{2}}} + \frac{c}{(\lambda_1 d_1)^{\frac{n}{2}}} + \frac{c}{(\lambda_2 d_2)^{\frac{n-2}{2}}} \frac{1}{(\lambda_1 d_1)^{\frac{n}{2}}} = o\left(\varepsilon_{12}^{\frac{n+1}{n-2}}\right). \end{aligned} \quad (4.20)$$

Now, we deal with the first integral in the right-hand side of (4.14). To this aim, we set  $\tilde{\delta}_2 := \delta_{(a_1, \lambda_2)}$ . We note that, using Lemma 4.5, we derive that

$$\delta_2 = \tilde{\delta}_2 + O\left(\lambda_2 |a_2 - a_1| \tilde{\delta}_2\right). \quad (4.21)$$

We now introduce the following sets:

$$\begin{aligned} A_1 &:= \{x \in \mathbb{R}^n : \delta_1 \leq \varepsilon_{12}^{1/6} \tilde{\delta}_2\} = \{x \in \mathbb{R}^n : |x - a_1| \geq \frac{\beta}{\sqrt{\lambda_1 \lambda_2}}\} \\ A_2 &:= \{x \in \mathbb{R}^n : \tilde{\delta}_2 \leq \varepsilon_{12}^{1/6} \delta_1\} = \{x \in \mathbb{R}^n : |x - a_1| \leq \frac{1}{\beta \sqrt{\lambda_1 \lambda_2}}\} \\ A_3 &:= \mathbb{R}^n \setminus (A_1 \cup A_2), \end{aligned}$$

where

$$\beta := \left( \lambda_1 - \varepsilon_{12}^{1/(3(n-2))} \lambda_2 \right)^{1/2} \left( \lambda_1 \varepsilon_{12}^{1/(3(n-2))} - \lambda_2 \right)^{-1/2} \sim \varepsilon_{12}^{-1/(3(n-2))} \quad \text{as } \varepsilon \rightarrow 0.$$

We will estimate the first integral in the right-hand side of (4.14) on each set  $A_i$  for  $i = 1, 2, 3$ . First, we write

$$\int_{A_1} |f|^{p-1-\varepsilon} f \psi_1 = - \int_{A_1} (\alpha_2 \delta_2)^{p-\varepsilon} \psi_1 + O \left( \int_{A_1} \delta_2^{p-1} \delta_1^{\frac{2n-2}{n-2}} |x - a_1| \right). \quad (4.22)$$

Recall that  $\varepsilon_{12}$  satisfies (4.6). Hence we get

$$\begin{aligned} \int_{A_1} \delta_2^{p-1} \delta_1^{\frac{2n-2}{n-2}} |x - a_1| &\leq \varepsilon_{12}^{\frac{n-4}{6(n-2)}} \int_{A_1} \delta_2^{\frac{n}{n-2}} \delta_1^{\frac{n+2}{n-2}} |x - a_1| \\ &\leq \varepsilon_{12}^{\frac{n-4}{6(n-2)}} \lambda_2^{\frac{n}{2}} \int_{A_1} \delta_1^{\frac{n+2}{n-2}} |x - a_1| = o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right), \end{aligned} \quad (4.23)$$

$$\begin{aligned} \int_{A_1} \delta_2^{p-\varepsilon} \psi_1 &= \int_{\mathbb{R}^n} \delta_2^{p-\varepsilon} \psi_1 - \int_{A_2 \cup A_3} \delta_2^{p-\varepsilon} \psi_1 \\ &= \int_{\mathbb{R}^n} \delta_2^p \psi_1 + O \left( \varepsilon \varepsilon_{12} (\log(\varepsilon_{12}^{-1}))^{\frac{n-2}{n}} \right) + O \left( \lambda_2 |a_1 - a_2| \int_{A_2 \cup A_3} \tilde{\delta}_2^p |\psi_1| \right), \end{aligned} \quad (4.24)$$

where we have used the evenness of  $\tilde{\delta}_2$ , the oddness of  $\psi_1$ , (2.8) and (4.21). Note that, as in (4.21), we get

$$\int_{A_2 \cup A_3} \tilde{\delta}_2^p |\psi_1| \leq \varepsilon_{12}^{\frac{-1}{6(n-2)}} \int_{A_2 \cup A_3} \tilde{\delta}_2^{\frac{n+1}{n-2}} |x - a_1| \delta_1^{\frac{n+1}{n-2}} \leq \varepsilon_{12}^{\frac{-1}{6(n-2)}} \varepsilon_{12}^{\frac{n+1}{n-2}} \log(\varepsilon_{12}^{-1}). \quad (4.25)$$

Combining (4.22)-(4.25), (ii) of Lemma 4.5 and using the estimate F11 of [1], we obtain

$$\int_{A_1} |f|^{p-1-\varepsilon} f \psi_1 = \alpha_2^{p-\varepsilon} \frac{c}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_1} + o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right). \quad (4.26)$$

Secondly, we write

$$\begin{aligned} \int_{A_2} |f|^{p-1-\varepsilon} f \psi_1 &= \int_{A_2} (\alpha_1 \delta_1)^{p-\varepsilon} \psi_1 - (p-\varepsilon) \int_{A_2} (\alpha_1 \delta_1)^{p-1-\varepsilon} \alpha_2 \delta_2 \psi_1 \\ &\quad + O \left( \int_{A_2} \delta_2^2 \delta_1^{\frac{6}{n-2}} |x - a_1| \right). \end{aligned} \quad (4.27)$$

Note that

$$\int_{A_2} \delta_2^2 \delta_1^{\frac{6}{n-2}} |x - a_1| \leq \varepsilon_{12}^{\frac{n-5}{6(n-2)}} \int_{A_2} \delta_2^{\frac{n+1}{n-2}} \delta_1^{\frac{n+1}{n-2}} |x - a_1| = o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right), \quad (4.28)$$

$$(p-\varepsilon) \int_{A_2} \delta_1^{p-1-\varepsilon} \delta_2 \psi_1 = p \int_{\mathbb{R}^n} \delta_1^{p-1} \delta_2 \psi_1 - p \int_{A_1 \cup A_3} \delta_1^{p-1} \delta_2 \psi_1 + O \left( \varepsilon \varepsilon_{12} (\log(\varepsilon_{12}^{-1}))^{\frac{n-2}{n}} \right). \quad (4.29)$$

Using the evenness of  $\tilde{\delta}_2$ , the oddness of  $\psi_1$  and (4.21), we get

$$\int_{A_1 \cup A_3} \delta_1^{p-1} \delta_2 \psi_1 = O \left( \lambda_2 |a_1 - a_2| \int_{A_1 \cup A_3} \delta_1^{p-1} \tilde{\delta}_2 |\psi_1| \right). \quad (4.30)$$

Arguing as in (4.23), we obtain

$$\int_{A_1 \cup A_3} \delta_1^{p-1} \delta_2 \psi_1 = o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right). \quad (4.31)$$

Thus, using the evenness of  $\delta_1$ , the oddness of  $\psi_1$ , the estimate F11 of [1] and combining (4.27)-(4.31), we derive that

$$\int_{A_2} |f|^{p-1-\varepsilon} f \psi_1 = -\alpha_1^{p-1-\varepsilon} \alpha_2 \frac{c}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_1} + o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right). \quad (4.32)$$

Lastly, denoting by  $\tilde{f} = \alpha_1 \delta_1 - \alpha_2 \tilde{\delta}_2$ , and using the evenness of  $\tilde{f}$ , we obtain

$$\int_{A_3} |f|^{p-1-\varepsilon} f \psi_1 = O \left( \int_{A_3} |\tilde{f}|^{p-1} |\delta_2 - \tilde{\delta}_2| |\psi_1| + \int_{2|\delta_2 - \tilde{\delta}_2| \geq |\tilde{f}|} |\delta_2 - \tilde{\delta}_2|^p |\psi_1| \right). \quad (4.33)$$

Note that the set  $\{x : 2|\delta_2 - \tilde{\delta}_2| \geq |\tilde{f}|\}$  implies that  $\delta_1 = \tilde{\delta}_2(1 + o(1))$ . Hence it is contained in  $F := \{x : |x - a_1| \geq c(\lambda_1 \lambda_2)^{-1/2}\}$ . Therefore, as in (4.23), we get

$$\int_{2|\delta_2 - \tilde{\delta}_2| \geq |\tilde{f}|} |\delta_2 - \tilde{\delta}_2|^p |\psi_1| \leq c(\lambda_2 |a_1 - a_2|)^{\frac{n+2}{n-2}} \int_F \tilde{\delta}_2^{\frac{n}{n-2}} |x - a_1| \delta_1^p = o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right). \quad (4.34)$$

Now, arguing as in (4.25), we find

$$\int_{A_3} |\tilde{f}|^{p-1} |\delta_2 - \tilde{\delta}_2| |\psi_1| \leq c(\lambda_2 |a_1 - a_2|) \int_{A_3} (\delta_1^{p-1} + \tilde{\delta}_2^{p-1}) \tilde{\delta}_2 |x - a_1| \delta_1^{\frac{n}{n-2}} = o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right). \quad (4.35)$$

Combining (4.33)-(4.35), we obtain

$$\int_{A_3} |f|^{p-1-\varepsilon} f \psi_1 = o \left( \varepsilon_{12}^{\frac{n+1}{n-2}} \right). \quad (4.36)$$

Clearly, our Proposition follows from (4.9)-(4.20), (4.26), (4.32) and (4.36).  $\square$

Now, we are ready to prove Theorem 1.8.

**Proof of Theorem 1.8** According to Remark 2.3 and Lemmas 4.2, 4.4 and 4.5, Claim (a) follows. Now we will prove (1.10) in the case where  $n \geq 6$ . On one hand, observe that Proposition 4.6 and Remark 4.3 imply

$$\left| \frac{1}{\lambda_2} \frac{\partial \varepsilon_{12}}{\partial a_2} \right| = \left| \frac{\lambda_1}{\lambda_2} \frac{1}{\lambda_1} \frac{\partial \varepsilon_{12}}{\partial a_2} \right| = o \left( \varepsilon_{12}^{\frac{n-1}{n-2}} \right). \quad (4.37)$$

On the other hand, Remark 4.3 implies that

$$\left| \frac{1}{\lambda_2} \frac{\partial \varepsilon_{12}}{\partial a_2} \right| = c\lambda_1 |a_1 - a_2| \varepsilon_{12}^{\frac{n}{n-2}} \geq c\sqrt{\lambda_1 \lambda_2} |a_1 - a_2| \varepsilon_{12}^{\frac{n-1}{n-2}}. \quad (4.38)$$

Clearly, from (4.37), (4.38) and Remark 2.3, we get the first claim of (1.10). Now, using (4.37), Lemma 4.4 and (2.11) for  $i = 2$ , we derive that

$$\frac{1}{\lambda_2^{n-1}} \frac{\partial H}{\partial a_2}(a_2, a_2) = o\left(\frac{1}{(\lambda_2 d_2)^{n-1}}\right).$$

This implies that

$$d_2 \not\rightarrow 0 \quad \text{and} \quad \frac{\partial H}{\partial a_2}(a_2, a_2) \rightarrow 0 \quad \text{as} \quad \varepsilon \rightarrow 0.$$

This completes the proof of (1.10).

It remains to prove Claim (b). Observe that, by Theorem 1.1 and assumption (1.9), we have

$$u_\varepsilon = P\delta_{(a_{\varepsilon,1}, \mu_{\varepsilon,1})} - P\delta_{(a_{\varepsilon,2}, \mu_{\varepsilon,2})} + v \quad \text{with} \quad \|v\| \rightarrow 0 \quad \text{and} \quad u_\varepsilon(a_{\varepsilon,1}) = \max |u_\varepsilon|. \quad (4.39)$$

First we claim that

$$\exists m > 0 \quad \text{s.t.} \quad h_\varepsilon := \max_{x \in \Omega} |x - a_{\varepsilon,1}|^{(n-2)/2} |u_\varepsilon(x)| \leq m. \quad (4.40)$$

In fact, if  $h_\varepsilon \rightarrow +\infty$ , then, using the method of R. Schoen [14], we can construct a concentration point  $b_\varepsilon$  with a concentration speed  $c|u_\varepsilon(b_\varepsilon)|^{2/(n-2)-\varepsilon/2}$  and the function  $u_\varepsilon$  becomes close to

$$P\delta_{(a_{\varepsilon,1}, \mu_{\varepsilon,1})} - P\delta_{(b_\varepsilon, c|u_\varepsilon(b_\varepsilon)|^{2/(n-2)-\varepsilon/2})}.$$

Observe that, since  $h_\varepsilon \rightarrow +\infty$ , we derive that

$$c|u_\varepsilon(b_\varepsilon)|^{2/(n-2)-\varepsilon/2} |b_\varepsilon - a_{\varepsilon,1}| \rightarrow \infty,$$

which contradicts the conclusions of Theorem 1.8. Hence our claim is proved.

Now, let

$$\Omega_1 := \Omega \setminus B(a_{\varepsilon,2}, \mu_{\varepsilon,2}^{\varepsilon(n-2)/4-1}),$$

and we introduce the following function

$$w_\varepsilon(X) := \frac{1}{\mu_{\varepsilon,2}^{(n-2)/2}} u_\varepsilon(a_{\varepsilon,2} + \mu_{\varepsilon,2}^{\varepsilon(n-2)/4-1} X) \quad \text{for} \quad X \in \Omega'_\varepsilon := \mu_{\varepsilon,2}^{1-\varepsilon(n-2)/4} (\Omega - a_{\varepsilon,2}).$$

The function  $w_\varepsilon$  satisfies

$$\begin{cases} -\Delta w_\varepsilon = |w_\varepsilon|^{2^*-2-\varepsilon} w_\varepsilon & \text{in } \Omega'_\varepsilon, \\ w_\varepsilon = 0 & \text{on } \partial\Omega'_\varepsilon. \end{cases}$$

Observe that, using (4.40) we derive that

$$|w_\varepsilon(x)| \leq m \quad \text{for each } x \in \tilde{\Omega}_1 := \Omega'_\varepsilon \setminus B(0, 1).$$

Hence  $w_\varepsilon$  converges in  $C_{loc}^2(\mathbb{R}^n \setminus B(0, 1))$  to a function  $w$  satisfying

$$-\Delta w = |w|^{2^*-2} w \quad \text{in } \mathbb{R}^n \setminus B(0, 1).$$



But from (4.39), we deduce that  $w$  has to be  $-\delta_{(0, \beta_n^{2/(n-2)})}$  (where  $\beta_n$  is defined in Theorem 1.1). Thus  $w < -c < 0$  in  $\partial B(0, 2)$  which implies that  $w_\varepsilon < 0$  in the same set. Hence  $u_\varepsilon < 0$  in  $\partial B_\varepsilon := \partial B(a_{\varepsilon, 2}, 2\mu_{\varepsilon, 2}^{\varepsilon(n-2)/4-1})$ . Now since  $\mu_{\varepsilon, 2}d(a_{\varepsilon, 2}, \partial\Omega) \rightarrow \infty$  (see Theorem 1.1) we derive that  $\overline{B_\varepsilon}$  is contained in a compact set of  $\Omega$ . Finally, using the fact that  $u_\varepsilon(a_{\varepsilon, 1}) > 0$ ,  $a_{\varepsilon, 1} \in B_\varepsilon$  (since  $\mu_{\varepsilon, 2}|a_{\varepsilon, 2} - a_{\varepsilon, 1}| \rightarrow 0$ ) and  $\Omega \setminus \{x : u_\varepsilon(x) = 0\}$  has exactly two connected components, we deduce that the nodal surface does not intersect the boundary of  $\Omega$ . Hence Claim (b) follows.

This completes the proof of our Theorem.  $\square$

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